

# Image Sensor Modeling: Color Measurement at Low Light Levels

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**Abstract.** The investigation of low light imaging is of high importance in the field of color science from different perspectives. One of the most important challenges that arises at low light levels is the issue of noise or, more generally speaking, low signal-to-noise ratio (SNR). In the present work, effects of different image sensor noises, such as photon noise, dark current noise, read noise, and quantization error, are investigated on low light color measurements. In this regard, a typical image sensor is modeled and employed for this study. A detailed model of noise is considered in the process of implementing the image sensor model to guarantee the precision of the results. Several experiments have been performed over the implemented framework and the results show the following: first, photon noise, read noise, and quantization error lead to uncertain measurements distributed around the noise free measurements and these noisy samples form an elliptical shape in the chromaticity diagram; second, even for an ideal image sensor, in very dark situations, stable measurement of color is impossible due to the physical limitation imposed by the fluctuations in photon emission rate; third, dark current noise reveals dynamic effects on color measurements by shifting their chromaticities towards the chromaticity of the camera black point; fourth, dark current dominates the other sensor noise types in the image sensor in terms of affecting measurements. Moreover, an SNR sensitivity analysis against the noise parameters is presented over different light intensities. © 2014 Society for Imaging Science and Technology.

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## INTRODUCTION

The human visual system is able to work under different lighting conditions. It is desirable to have imaging devices, such as cameras, that are able to operate in similar light levels. The ability of the human visual system to work even under low light situations leads to the importance of studying low light levels. However, most of the theories, measures, models, and methods in color science are developed for high intensities.<sup>1</sup> These theories, measures, methods, and models cannot be used for low light situations, since they fail to comply with the necessary conditions for which they are feasible. For instance, a color difference formula that is derived for photopic conditions (i.e., luminance levels greater than 5 cd/m<sup>2</sup>) cannot be leveraged in evaluating techniques developed for assessing dark images.<sup>2</sup> This issue implies the importance of investigating low light conditions. Moreover, the addressing of low light or, more generally speaking, low signal-to-noise level situations has a wide range of

applications in photography,<sup>3</sup> designing biosensors,<sup>4</sup> image processing,<sup>5</sup> machine vision, and color science.<sup>6</sup>

For both imaging devices and the human visual system, as the light level goes down, the effect of noise becomes more significant.<sup>6</sup> In this situation, an imaging device will acquire a noisy signal with a low value of signal-to-noise ratio (SNR). In human vision, lower signal-to-noise level results in changes in the appearance of measured colors. Several works have discussed the impact of light level on human color perception.<sup>7-10</sup> It is unanimously accepted that reducing the light level gives rise to color shifts, and this effect is mostly attributed to the rod intrusion into the mesopic vision (dim light situation in which both rods and cones contribute to vision).<sup>11,12</sup> However, the issue of noise at low light levels is still an open problem for artificial image sensors.

To the best of our knowledge, the effect of noise at low light levels on the color measurements of imaging devices has not been addressed yet. One of the most recent works concerning this topic is the work performed by Kirk and O'Brien, proposing a tone mapping approach to convert high dynamic low light images to a perceptually closer result to the human mesopic vision experience.<sup>13</sup> However, the authors did not take into account any noise type in their mesopic color appearance modeling and left it as a future work.

Our article is concerned with modeling the performance of color image sensors under low signal-to-noise ratios. Our methodology involves tracking photons in the imaging sensor pipeline from the emission to the detection and recording level. In this regard, physical rules governing photon emission are employed to estimate the low light version of quantities describing the light coming to the imaging device; then, an image sensor model is implemented and leveraged to study the luminance and noise induced effects on the sensor color measurements. Camera or image sensor models have been presented in different works.<sup>14,15</sup> The rationale behind modeling digital camera imaging systems is, first, to reconstruct hyperspectral images taken by spectrometers, or to be used in computer graphics applications<sup>16</sup>, or, second, to evaluate the camera design and output image quality, or to optimize the performance of the camera in terms of some adjustable parameters (e.g., exposure time or ISO setting).<sup>17,18</sup>

In terms of application, the results of this study can be utilized in developing low light image quality measures, introducing efficient denoising algorithms, developing realistic color noise perception models,<sup>19</sup> addressing low

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92 signal-to-noise levels in digital cameras, and developing an  
 93 automatic labeling system for micro-array sensor biochips.<sup>20</sup>

94 The remainder of this article is organized as follows.  
 95 The background for light emission is presented in the  
 96 second section followed by introduction of the image sensor  
 97 modeling in the third section. Experiments, results, and  
 98 discussion are reported in the fourth section. Finally, the fifth  
 99 section concludes the article.

## 100 BACKGROUND

101 Light consists of energy packets called *photons*. Photons are  
 102 characterized by their frequency and polarization state.<sup>21</sup>  
 103 Each photon carries an amount of energy determined  
 104 by its frequency. This energy is equal to  $\{hc/\lambda\}$ , where  
 105  $h = 6.626,176 \times 10^{-34}$  (J s) is the Planck constant,  $c =$   
 106  $2.997,925 \times 10^8$  (m/s) represents the light speed, and  
 107  $\lambda$  denotes the wavelength (inverse of frequency) of the  
 108 photon. Photon emission from a light source follows a  
 109 Poisson distribution. For a monochromatic light source of  
 110 a particular wavelength  $\lambda_0$  and known average number of  
 111 emitted photons per second  $g$ , the probability of emitting  $n$   
 112 photons per unit of time can be obtained as follows<sup>22</sup>:

$$113 \quad P(g, n) = \frac{g^n e^{-g}}{n!}. \quad (1)$$

114 Given the spectral radiance,  $L(\lambda)$ , the average emitted  
 115 number of photons, per unit time, per unit area, per unit  
 116 steradian, for a central wavelength  $\lambda_0$  can be obtained by  
 117 calculating the following integral over an infinitely small  
 118 range of  $[\lambda_0 - \delta/2, \lambda_0 + \delta/2]$ :

$$119 \quad g(\lambda_0) = \frac{1}{hc} \int_{\lambda_0 - \delta/2}^{\lambda_0 + \delta/2} \lambda L(\lambda) d\lambda. \quad (2)$$

120 The wavelength range of the spectrum,  $[\lambda_{\min}, \lambda_{\max}]$ , can  
 121 be discretized into  $N$  intervals of the length  $\delta$  such that  
 122  $\{\lambda_{\max} - \lambda_{\min} = N\delta\}$ . Hence,  $g(\lambda_i)$  of the  $i$ th wavelength bin  
 123 can be approximated as

$$124 \quad g(\lambda_i) = \frac{1}{hc} \int_{\lambda_i - \delta/2}^{\lambda_i + \delta/2} \lambda L(\lambda) d\lambda \approx \frac{\lambda_i L(\lambda_i) \delta}{hc}. \quad (3)$$

125 Let  $L(\lambda)$  represent the high intensity radiance of a light. Our  
 126 goal is to derive an estimation of this spectral radiance at an  
 127 arbitrary lower intensity. The high intensity spectral radiance  
 128 is the most complete description of the light and this quantity  
 129 at any lower intensity can be predicted from the given high  
 130 intensity spectrum as described in the following.

131 The Poisson distribution,  $\text{Pois}(g(\lambda_i))$ , corresponding to  
 132 each bin of the high intensity spectral radiance is fully  
 133 characterized by knowing the  $g(\lambda_i)$  values. We define the  
 134 intensity factor  $F \leq 1$ , which is a multiplier to change the  
 135 light level. The estimated spectral radiance after applying  
 136 the intensity factor  $F$  can be obtained by drawing samples,  
 137  $\{\tilde{G}_F(\lambda_i)\}_1^N$ , from  $\{\text{Pois}(F \times g(\lambda_i))\}_1^N$  distributions. Hence,  
 138 the estimated spectral radiance,  $\tilde{L}_F(\lambda)$ , for the intensity factor

$F$  and central wavelength  $\lambda_i$  is given by

$$139 \quad \tilde{L}_F(\lambda_i) = \frac{\tilde{G}_F(\lambda_i) \times hc}{\lambda_i \delta}. \quad (4) \quad 140$$

141 By taking this approach, we can establish the effect of  
 142 shot noise on estimates of low light spectral radiances. It  
 143 is worth mentioning that  $\tilde{L}_{FN}(\alpha, \beta, \lambda)$ , which denotes the  
 144 quantal number of photons falling on the location  $(\alpha, \beta)$  of  
 145 the image sensor in photons/sec/m<sup>2</sup>/sr/nm, can be obtained  
 146 from the radiance quantity,  $\tilde{L}_F(\alpha, \beta, \lambda)$ , as

$$147 \quad \tilde{L}_{FN}(\alpha, \beta, \lambda) = \frac{\tilde{L}_F(\alpha, \beta, \lambda) \times \lambda}{hc}. \quad (5) \quad 148$$

## 148 IMAGE SENSOR MODELING

149 A typical digital camera is comprised of the following  
 150 elements: an optical system, an image sensor, and an image  
 151 processor.<sup>16</sup> The focus of this section is on modeling  
 152 and simulating the image sensor of a digital camera. We  
 153 consider the image formation model, noise model, and  
 154 analog-to-digital converter (ADC) components in the image  
 155 sensor model. Figure 1 shows a diagram of the image sensor  
 156 model, which is a modified version of the model introduced  
 157 in Ref. 18.

158 When the shutter of a camera opens, a stream of photons  
 159 enters the camera and falls on the image sensor. A color  
 160 image sensor consists of three sensor types, which usually  
 161 are referred to as R, G, and B sensors. The exposure setting  
 162 determines the number of photons captured by the sensors.  
 163 Each sensor type has a specific spectral quantum efficiency  
 164 (i.e., the proportion of electrons generated as a result of  
 165 photon catches for an area of 1 (m<sup>2</sup>) that subtends 1 (sr)). A  
 166 pixel of an image sensor consists of a photodetector, a color  
 167 filter, and a readout circuit. The rain of photons hitting the  
 168 photodetector produces a photocurrent. This photocurrent  
 169 together with the photodetector dark current, which will be  
 170 explained later, is accumulated during the integration time  
 171 as far as the sensor capacity allows. The maximum sensor  
 172 charge capacity is known as full-well capacity and determines  
 173 the level of saturation for each sensor. When the integration  
 174 time is over, the readout circuit is responsible for measuring  
 175 the voltage produced in the pixels. This process is prone  
 176 to noise, known as the readout noise. The structure of the  
 177 readout circuit is the main difference between CCD and  
 178 CMOS type image sensors.

### 179 Noise Model

180 Noise can be defined as any unwanted event that hampers  
 181 the image quality. In our simulation framework, we assume  
 182 an additive model for the noise and the following noise  
 183 types are considered as the most significant sources of noise  
 184 underlying the image distortion.

#### 185 Photon Shot Noise

186 Photon shot noise can be defined as the variation in the  
 187 number of photons emitted from the light source and,

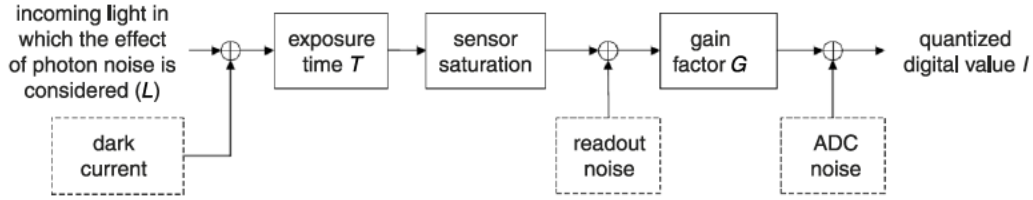


Figure 1. Image sensor prototype for a single channel.

188 consequently, the number of photons detected in the image  
 189 sensor at different times. This phenomenon is rooted in the  
 190 probabilistic nature of photon emission, as explained in the  
 191 second section.

### 192 Dark Current Noise

193 The current produced inside the image sensor in the absence  
 194 of light is referred to as the *dark current noise*. This current  
 195 is not generated as a result of photogeneration but as  
 196 a result of the impurities existing in the silicon wafer.<sup>23</sup>  
 197 Dark current noise is also known as thermal noise and  
 198 ambient temperature has a large influence on its amplitude.  
 199 Dark current introduces shot noise to the measurement<sup>23</sup>  
 200 and can be modeled by a Poisson distribution with a  
 201 variance of  $(\sigma_{\text{dark}}^{\kappa})^2$  for the  $\kappa$  sensor type. Since the variance  
 202 of a Poisson distribution is equal to its expected value,  
 203 the parameter  $(\sigma_{\text{dark}}^{\kappa})^2$  represents the average number of  
 204 generated electrons as a result of dark current for each pixel  
 205 per unit time.

$$206 \quad N_{\text{dark}}^{\kappa}(\alpha, \beta) \sim \text{Pois}((\sigma_{\text{dark}}^{\kappa})^2). \quad (6)$$

### 207 Read Noise

208 Read noise refers to the noise in the readout circuit, caused by  
 209 an on-chip amplifier, and can be modeled as a white Gaussian  
 210 distribution with standard deviation  $\sigma_{\text{read}}$ .<sup>14</sup> Readout noise  
 211 limits the dynamic range of image sensors.

$$212 \quad N_{\text{read}} \sim N(0, \sigma_{\text{read}}). \quad (7)$$

### 213 Quantization Noise

214 In the last step of generating the digital image in the image  
 215 sensor prototype, the amplified voltage should be quantized  
 216 into digital values. Quantization error introduced in this  
 217 step is known as quantization noise and represented as  $\sigma_{\text{adc}}$ .  
 218 The noise induced by the amplifier of the analog-to-digital  
 219 conversion unit (ADC) is considered negligible.

### 220 Pixel Measurement Model

221 The voltage produced by an image sensor can be obtained by  
 222 the following formula:

$$223 \quad V^{\kappa}(\alpha, \beta) = G_{\text{Ve}^-} \times f_{\text{sat}} \left( T \times \int_{\lambda_{\min}}^{\lambda_{\max}} \tilde{L}_{FN}(\alpha, \beta, \lambda) Q_e^{\kappa}(\lambda) d\lambda \right. \\ 224 \quad \left. + T \times N_{\text{dark}}^{\kappa}(\alpha, \beta) \right), \quad (8)$$

225 where  $\kappa \in \{R, G, B\}$ ,  $T$  indicates the exposure time in s,  
 226  $G_{\text{Ve}^-}$  is the conversion gain in  $\text{V}/e^-$ ,  $\tilde{L}_{FN}(\alpha, \beta, \lambda)$  represents  
 227 the number of incident photons at the location  $(\alpha, \beta)$   
 228 of the image sensor obtained from the spectral radiance  $\tilde{L}_F$   
 229 at intensity factor  $F$  in  $\text{photons}/\text{s}/\text{m}^2/\text{sr}/\text{nm}$ ,  $Q_e^{\kappa}(\lambda)$  is the  
 230 quantum efficiency of the  $\kappa$  sensor, in  $e^- \text{ m}^2 \text{ sr}/\text{photons}$ ,  
 231  $N_{\text{dark}}^{\kappa}(\alpha, \beta)$  represents the number of electrons generated as  
 232 a result of dark noise in the  $\kappa$  channel for the pixel  $(\alpha, \beta)$ , and  
 233  $f_{\text{sat}}(\cdot)$  indicates the saturation function of the sensor.

234 The quantum efficiency curve for the  $\kappa$  sensor type  
 235 is defined as the proportion of the number of electrons  
 236 generated by the sensor,  $N_e^{\kappa}$ , to the number of incident  
 237 photons with the wavelength  $\lambda$ ,  $N_{\text{ph}}^{\kappa}$ .<sup>24</sup>

$$238 \quad Q_e^{\kappa}(\lambda) = \frac{N_e}{N_{\text{ph}}^{\kappa}(\lambda)}. \quad (9) \quad 239$$

239 The voltage measured by the readout circuit is given by

$$240 \quad \tilde{V}^{\kappa}(\alpha, \beta) = V^{\kappa}(\alpha, \beta) + N_{\text{read}}(\alpha, \beta). \quad (10)$$

241 The raw output image of the camera can be obtained after  
 242 applying the gain factor and then the quantization process as  
 243 follows:

$$244 \quad I^{\kappa}(\alpha, \beta) = [G \times \tilde{V}^{\kappa}(\alpha, \beta)]_{n_b}. \quad (11)$$

245 In the above equation,  $[\cdot]_{n_b}$  represents the “ $n_b$ -bit” quanti-  
 246 zation operation that outputs the integer part of the given  
 247 operand,  $G \times \tilde{V}^{\kappa}(\alpha, \beta)$ , in the range of  $[0, 2^{n_b} - 1]$ . Hence,  
 248 the quantization noise of the  $\kappa$ th channel at the location  
 249  $(\alpha, \beta)$  of the image is given by

$$250 \quad \sigma_{\text{ADC}}(\alpha, \beta)^{\kappa} = I^{\kappa}(\alpha, \beta) - G \times \tilde{V}^{\kappa}(\alpha, \beta). \quad (12)$$

251 Finally, the signal-to-noise ratio can be defined as the  
 252 ratio of the nonsaturated output of the noise free signal to  
 253 the variance of the noise. The total variance of noise for  
 254 each sensor type at each pixel location can be estimated as  
 255 follows<sup>25</sup>:

$$256 \quad \text{Var}^{\kappa}(\alpha, \beta) = V^{\kappa}(\alpha, \beta) \times G^2 + \sigma_{\text{read}}^2 \times G^2 + (\sigma_{\text{ADC}}^{\kappa}(\alpha, \beta))^2. \quad (13)$$

257 For nonsaturated pixels in the image, the SNR value of each  
 258 channel can be obtained by the following formula<sup>18</sup>:

$$259 \quad \text{SNR}^{\kappa}(\alpha, \beta) \\ 260 \quad = \frac{\left[ G \times G_{\text{Ve}^-} \times T \times \int_{\lambda_{\min}}^{\lambda_{\max}} \tilde{L}_{FN}(\alpha, \beta, \lambda) Q_e^{\kappa}(\lambda) d\lambda \right]_{n_b}^2}{V^{\kappa}(\alpha, \beta) \times G^2 + \sigma_{\text{read}}^2 \times G^2 + \sigma_{\text{ADC}}^2}. \quad (14) \quad 261$$

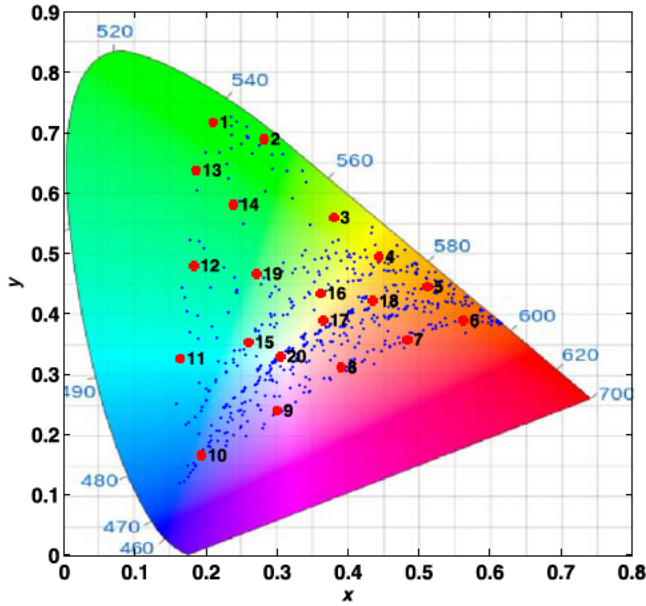


Figure 2. The chromaticity values spanned by the RGB598 spectral database are indicated by blue dots. The selected data points for the experiment are marked as red asterisks with designated numeric indices.

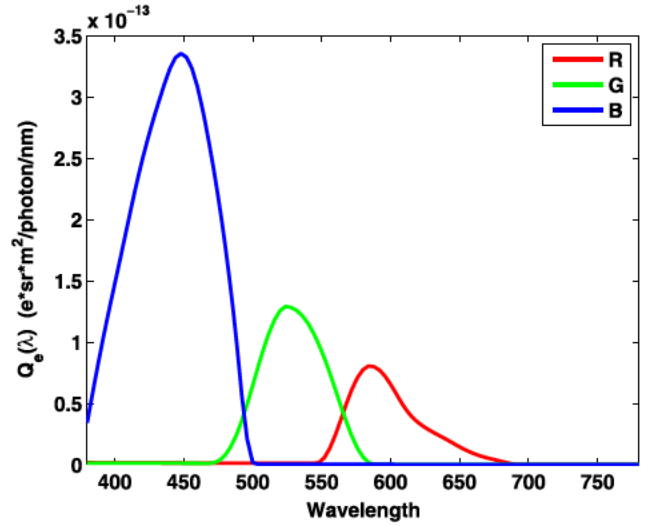


Figure 3. The quantum efficiency curves of image sensors in  $e^- \text{ sr m}^2 / \text{photon / nm}$ .

## 262 EXPERIMENTS, RESULTS AND DISCUSSION

263 We designed a set of experiments intended to investigate the  
 264 effects of different noise types on the color measurements  
 265 of image sensors. The experiments were performed using  
 266 the spectral radiances selected from “A Data Set for Color  
 267 Research,” prepared by Barnard et al.<sup>26</sup> The dataset contains  
 268 the spectral sensor sensitivity curves of the Sony DXC-930  
 269 three-chip CCD video camera and the spectra of 23 of  
 270 the Macbeth color patches illuminated by 26 different light  
 271 sources. The Sony DXC-930 sensor sensitivity curves are  
 272 used in the image sensor simulation phase of this work  
 273 and the spectra, which we refer to as the RGB598 spectral  
 274 database, are leveraged for our experiments and evaluation  
 275 of the simulation. The sensor quantum efficiency curves are  
 276 shown in Figure 3. Each spectrum is sampled in 4 nm steps  
 277 from 380 to 780. Details about this database can be found in  
 278 Ref. 26. The chromaticities spanned by the 598 spectra of this  
 279 database are shown in the chromaticity diagram in Figure 2,  
 280 and, among these points, 20 points are selected for the sake  
 281 of our experiments. First, each spectral radiance is scaled  
 282 to have a spectrum with a luminance value of 100 ( $\text{cd/m}^2$ );  
 283 then, to obtain a lower luminance value, the spectrum is  
 284 used as  $L(\lambda)$  in Eqs. (3)–(5) to estimate the corresponding  
 285 low intensity spectral radiance  $\tilde{L}_F$  at the intensity factor  $F$ . It  
 286 is worth mentioning that since the luminance of the scaled  
 287 spectral radiances are set to 100 at the intensity factor  $F = 1$ ,  
 288 the approximate luminance value of  $\tilde{L}_F$  can be obtained as  
 289  $F \times 100$ .

290 For each data point, the raw output of the image sensor  
 291 is generated from the modeled framework at a specific  
 292 condition defined for each experiment. The parameters  
 293 selected for the image sensor model at a temperature of 20°C  
 294 are listed in Table I. The camera black RGB for Sony DXC-930

Table I. Parameters of the model at 20°C.

| Sensor parameter  | Parameter value |
|---|-----------------|
| $G_{Ve^-}$ ( $V/e^-$ )  | 0.0002          |
| $(\sigma_{\text{dark}}^{\kappa})^2$ ( $e^- / \text{pixel} / \text{s}$ ) | [195, 230, 218] |
| $\sigma_{\text{read}}$ ( $e^-$ )  | 4               |
| $G$   | 141.67          |
| Full-well capacity ( $e^-$ )  | 9000            |
| $T$ (s)   | 1               |
| $\eta_b$  | 8               |

is provided in the RGB598 database and this value is scaled to  
 295 obtain the variance of dark noise  $(\sigma_{\text{dark}}^{\kappa})^2$ . Full-well capacity,  
 296 read noise standard deviation ( $\sigma_{\text{read}}$ ), and the conversion  
 297 gain ( $G_{Ve^-}$ ) are selected from Ref. 16. Based on these selected  
 298 values, the parameter  $G$  is determined such that the output of  
 299 the sensor best fits the empirical measurements given in the  
 300 RGB598 database.

To account for uncertainties imposed by noise, 200  
 302 measurements are recorded for each sample in each trial. The  
 303 measured samples ( $I$ ) are converted to the XYZ space ( $I_{XYZ}$ ),  
 304 and then to the  $xy$ -chromaticity space. This transformation  
 305 is given by  
 306

$$I_{XYZ} = M \times I, \quad (15)$$

$$M = (T_{XYZ} \times T_{XYZ}^t) \times (C \times T_{XYZ}^t)^{-1}, \quad 307$$

where  $T_{XYZ}$  and  $C$  are  $(3 \times N)$  matrices representing the  
 308 XYZ color matching function and the camera sensitivity  
 309 curves respectively. The camera sensitivity curves can be  
 310 obtained from the quantum efficiency function,  $Q_e^{\kappa}(\lambda_i)$ , as  
 311 follows:  
 312

$$C^{\kappa}(\lambda) = G_{Ve^-} \times G \times Q_e^{\kappa}(\lambda),$$

$$Q_e^{\kappa}(\lambda_i) = \frac{hc}{\lambda_i} \times Q_e^{\kappa}(\lambda_i) \quad \kappa \in \{R, G, B\}, \quad i \in \{1, 2, \dots, N\}. \quad (16) \quad 313$$

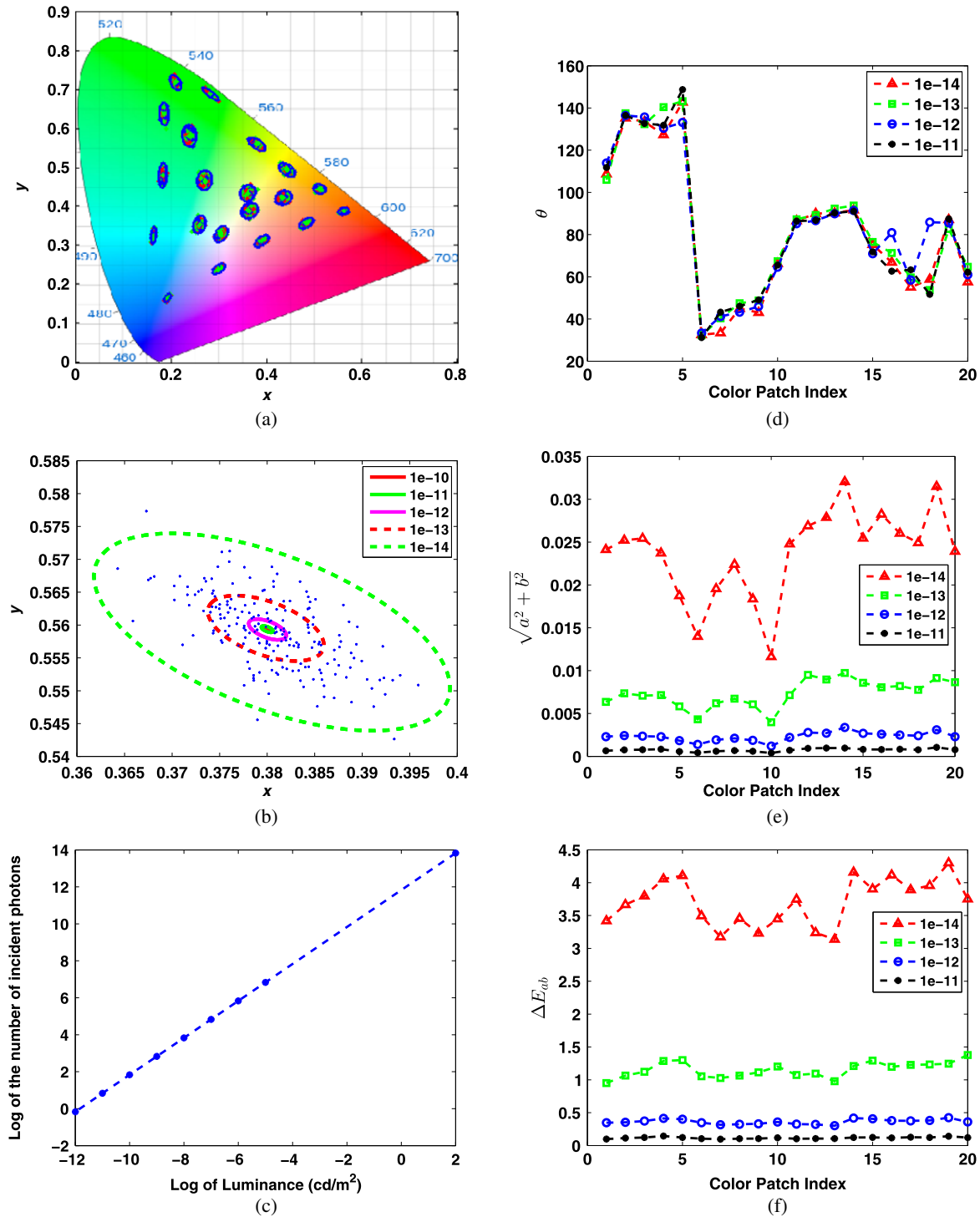


Figure 4. The results of scenario I performed over the RGB598 database. (a) The generated samples for each selected data point of the RGB598 database together with the fitted ellipses to those samples are plotted. (b) The generated samples and the fitted ellipses for different intensity factors for data point number 3. (c) The log number of incident photons at different luminance levels for color patch number 3 is depicted. (d) The estimated inclination angles of ellipses obtained from the PCA algorithm for different intensity factors are shown for all color patches. (e) The sizes of fitted ellipses corresponding to different intensity factors for all selected color patches are compared. (f) The average of  $\Delta E_{ab}$  values over the samples of each intensity factor.

314 Here a question may arise, which is “is it correct to use  
 315 CIE photopic colorimetry in a low light mesopic range?”  
 316 There are two reasons behind the choice of CIE XYZ  
 317 color matching functions: first, CIE photopic colorimetry is  
 318 commonly used in cameras in the process of creating the

output image; second, to the best of our knowledge, current  
 camera technology does not use any known color space or  
 model specific for mesopic conditions. A CIE system for  
 mesopic photometry was proposed in Ref. 27; however, the  
 use of it is still not widespread.

324 The experiments were carried out over three scenar- 378  
 325 ios and followed by an SNR sensitivity analysis. Before 379  
 326 explaining the results, we state the main assumptions and 380  
 327 considerations of this work as follows. 381

- 328 (1) The temperature is assumed constant, hence the dark 382  
 329 noise parameters are fixed during the experiments. 383
- 330 (2) The noise model is additive in the image sensor 384  
 331 simulation framework. 385
- 332 (3) The image sensor linearly responds to light intensity 386  
 333 variations before its saturation limit. Sensor linearity is 387  
 334 discussed in Ref. 28 more profoundly. In Ref. 28, Barnard 388  
 335 and Funt mention that “The Sony DXC-930 camera that 389  
 336 we used for our experiments is quite linear for most of its 390  
 337 range, provided it is used with gamma disabled.” 391
- 338 (4) Raw output images are considered for our analysis. 392
- 339 (5) *Reset noise*, *photodetector response nonuniformity* 393  
 340 (PRNU), and *dark signal nonuniformity* (DSNU) are 394  
 341 not incorporated in our modeling, and, for our research, 395  
 342 we assume that their impact on the introduced model 396  
 343 is negligible. For further details the reader can refer to 397  
 344 Ref. 25. 398

### 345 **Scenario I: Ideal Image Sensor and Light Intensity**

346 In the first scenario, we consider the case where there is no 401  
 347 noise corrupting the output image, we have a perfect image 402  
 348 sensor able to detect single photon events, and the sensor can 403  
 349 respond without saturation. We want to investigate at which 404  
 350 luminance value the photon noise becomes significant. In 405  
 351 this regard, twenty data points shown in Fig. 2 are considered 406  
 352 for this experiment. The log of the intensity factor is set to val- 407  
 353 ues  $\log(F) \in \{0, -7, -8, -9, -10, -11, -12, -13, -14\}$ . 408  
 354 The results of the experiment are shown in Figure 4. Fig. 4(a) 409  
 355 indicates that the generated samples form an elliptic shape in 410  
 356 the chromaticity diagram. The principal component analysis 411  
 357 (PCA) algorithm is used to find a fitted ellipse for the 412  
 358 generated samples of each data point.<sup>29</sup> 413

359 The generated samples and the fitted ellipses of the third 414  
 360 data point for different intensity factors, and the number of 415  
 361 incident photons on the image sensor for various luminance 416  
 362 values are plotted in Figs. 4(b) and 4(c) respectively. In 417  
 363 Fig. 4(b), the distance between consecutive ellipses grows 418  
 364 exponentially as the light intensity decreases. Figs. 4(d) 419  
 365 and 4(e) show the inclination angle and size of the fitted 420  
 366 ellipses for some intensity factors. The size of each ellipse 421  
 367 is approximated as  $\sqrt{a^2 + b^2}$ , where  $a$  and  $b$  represent 422  
 368 the sizes of the semi-major and semi-minor axes of the 423  
 369 ellipse. The inclination angle represents the angle between 424  
 370 the semi-major axis and the  $x$ -axis of the  $xy$ -chromaticity 425  
 371 space. The results indicate that the inclination angles, with 426  
 372 a good approximation, are independent of the intensity 427  
 373 level; however, the size of the ellipses inversely changes with 428  
 374 intensity, proving that even if we had an ideal image sensor 429  
 375 with no internal noise, we still had to deal with the photon 430  
 376 noise and uncertainties imposed by physical limitations. It 431  
 377 should be borne in mind that the photon noise becomes

noticeable at very low light levels, when the number of 378  
 incident photons is less than 100. Since distances in the 379  
 chromaticity diagram do not correspond to the human visual 380  
 system color discriminability,  $\Delta E_{ab}$  is used as a figure of 381  
 merit to show to what extent the effect of noise on color 382  
 measurement at different intensities would be noticeable to 383  
 a human observer from trial to trial. In this regard, for each 384  
 data point, the  $\Delta E_{ab}$  measure is derived as follows: 385

- 386 (1) the standard D65 illuminant is assumed as a white 387  
 388 reference for the calculations at the luminance of 389  
 100  $\text{cd/m}^2$  (the  $Y$  value of the reference white is kept 390  
 constant during the entire experiment); 391
- 392 (2) the  $XYZ$  values of each sample are scaled to equalize the 393  
 394  $Y$  value of the sample and that of the standard illuminant 395  
 to be able to compare the color coordinates of the low 396  
 intensity samples ( $F < 1$ ) and the high intensity sample 397  
 generated at ( $F = 1$ ); 398
- 399 (3) CIELab coordinates of each sample are obtained; 400
- 401 (4)  $\Delta E_{ab}$  is calculated between each sample and the average 402  
 403 color coordinates of corresponding high intensity sam- 404  
 ples; 405
- 406 (5) the average of  $\Delta E_{ab}$  values over the samples of each 407  
 408 intensity factor is reported. 409

The result of  $\Delta E_{ab}$  calculation is shown in Fig. 4(f), 410  
 indicating that as the light level falls off, the color change 411  
 between different trials of each data point becomes more 412  
 noticeable. 413

### 414 **Scenario II: Effects of Dark Current on Image Sensor** 415 **Responses at Low Light Intensity**

416 It is shown in the first scenario that photon noise may 417  
 418 bring about uncertainties in the measurements at very 419  
 420 low light levels when the image sensor is deemed ideal 421  
 422 and no other noises may disturb the measurement. In 423  
 424 this subsection, the effect of dark current is examined 425  
 426 separately from the other intrinsic noise types, when only 427  
 428 photon noise and dark current are affecting the image 429  
 430 sensor, and the sensor saturation function is not considered 431  
 in the sensor model. The intensity factor is set to  $F \in$   
 $\{1, 0.5, 0.1, 0.05, 0.01, 0.005, 0.001\}$  during each trial of 416  
 the experiment. For the sake of this experiment, only the 417  
 boundary data points (indices 1–13) from the initial 20 data 418  
 points are picked to make the resulting figures clearer. 419

420 The results are shown in Figure 5, indicating that the 421  
 422 dark noise may cause much more significant effects on 423  
 424 the color measurement at lower intensities than the photon 425  
 426 noise. The result is that the dark noise pushes the low 427  
 428 intensity measurements towards the average chromaticity 429  
 430 of the image sensor’s black point. In comparison to the 431  
 photon noise, which became noticeable at intensity factors 426  
 of the order of  $10^{-13}$  and lower, the dark current noise 427  
 effect becomes visible for  $F \leq 0.1$ . This issue indicates the 428  
 greater effects of dark noise compared with photon noise 429  
 in hampering the quality of measurements. The angle of 430  
 inclination of the ellipses,  $\theta$ , induced by the dark noise is 431

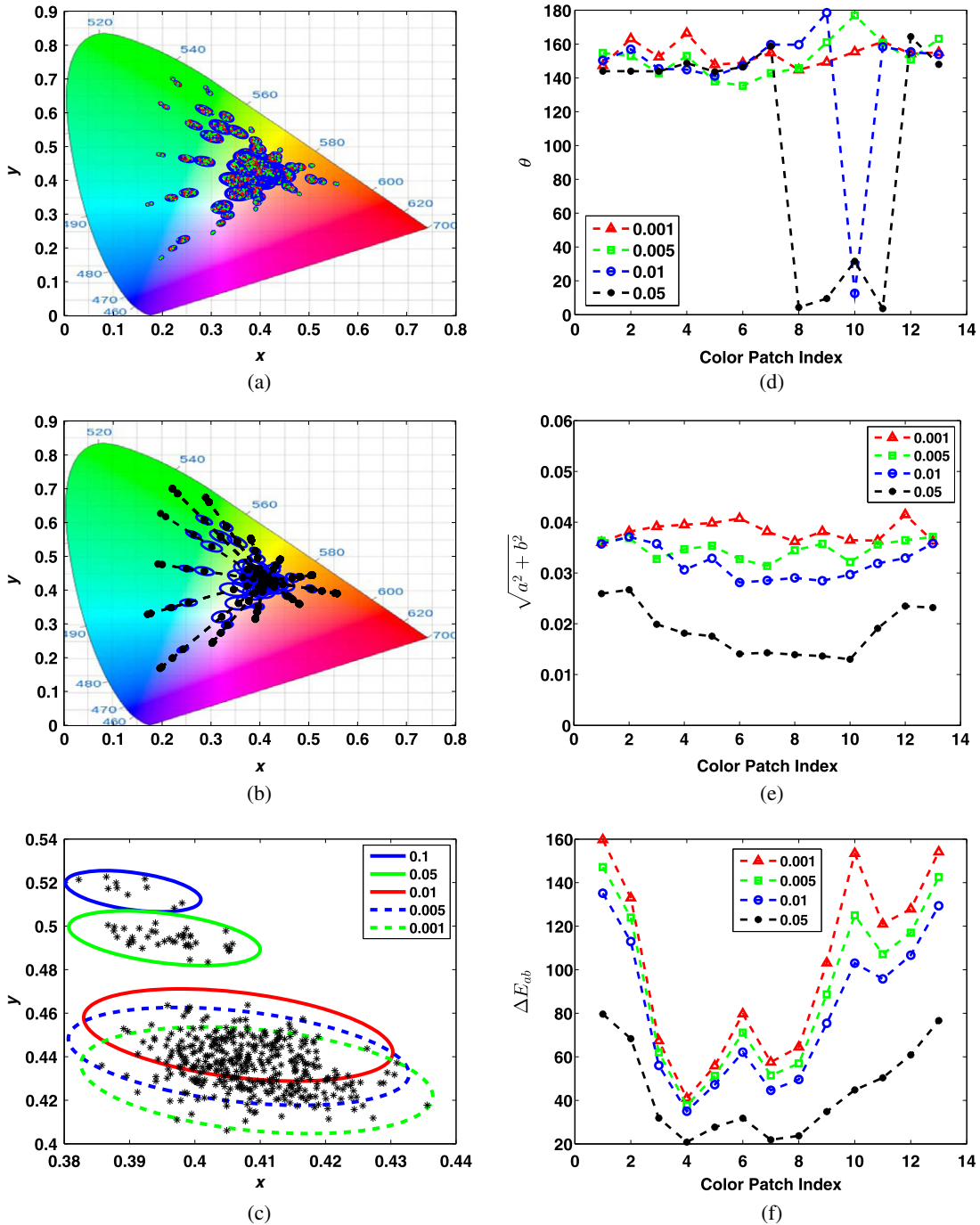


Figure 5. The results of scenario II performed over the RGB598 database when only photon noise and dark noise are taken into account in the image formation model. (a) Samples drawn for each selected data point of the RGB598 database and the ellipses fitted to the samples are plotted. (b) The subfigure in part (a) is regenerated after removing the samples and specifying the centers of the ellipses together with the line of movement of each data point with the light level. (c) The result of subfigure (a) is magnified for data point number 3. (d) The estimated inclination angles of ellipses obtained from the PCA algorithm for different intensity factors are shown for all color patches. (e) The sizes of fitted ellipses corresponding to different intensity factors for all selected color patches are compared. (f) The average of  $\Delta E_{ab}$  values over the samples of each intensity factor.

432 totally different from that of the photon noise. The ellipses  
 433 are aligned more horizontally for low intensities, and their  
 434 angles of inclination are more separated from each other  
 435 in different intensity factors than the results of scenario I.  
 436 Another interesting point is the opposite behavior of the  
 437 ellipse size variations as a function of the color patch index

in different light intensities. In scenario I, the size of the  
 ellipses is more uniform for lower intensity factors than for  
 higher values of  $F$ ; however, in scenario II, the opposite of  
 this pattern is exhibited, as seen in Fig. 5(e), where the size of  
 lower intensity ellipses is more uniform than high intensity  
 values.

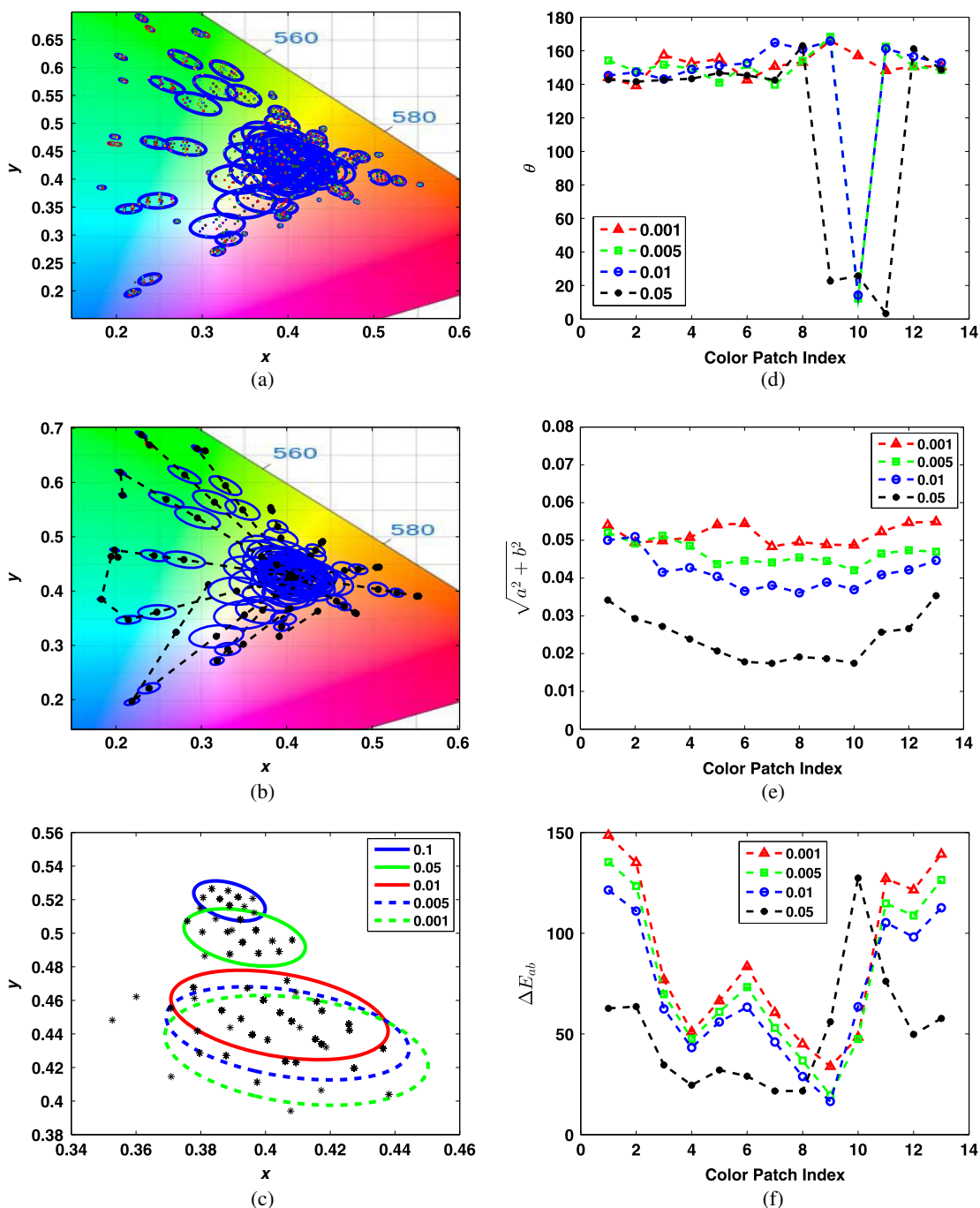


Figure 6. The results of scenario III performed over the RGB598 database when only photon noise and dark noise are taken into account in the image formation model. (a) Samples drawn for each selected data point of the RGB598 database and the ellipses fitted to the samples are plotted. (b) The subfigure in part (a) is regenerated after removing the samples and specifying the centers of the ellipses together with the line of movement of each data point with the light level. (c) The result of subfigure (a) is magnified for data point number 3. (d) The estimated inclination angles of ellipses obtained from the PCA algorithm for different intensity factors are shown for all color patches. (e) The sizes of fitted ellipses corresponding to different intensity factors for all selected color patches are compared. (f) The average of  $\Delta E_{ab}$  values over the samples of each intensity factor.

#### 444 Scenario III: Real Image Sensor Simulation

445 A similar scenario to scenario II is obtained with all  
 446 noise types and the saturation function being active. In  
 447 this experiment, only data points 1–13 are used. Figure 6  
 448 depicts the results. In Figs. 6(a) and 6(b), some data  
 449 points make the sensor saturated at high intensity factors.  
 450 Nonlinear effects imposed by these saturated samples are

explicitly revealed in Fig. 6(b). Moreover, the quantization 451  
 level in the model leads to sparse samples in the chromaticity 452  
 diagram, since it is not possible to have all chromaticity 453  
 values in the output of the image sensor. Aside from this, 454  
 the pattern of results of this scenario resembles that of 455  
 scenario II, implying the dominant influence of dark noise in 456  
 low light levels. 457



### 458 SNR Sensitivity Analysis

459 In this subsection, an analysis of the sensitivity of the SNR  
 460 value (given in Eq. (14)) to the parameters of dark current  
 461 and read noise and including or excluding the quantization  
 462 noise is presented. In this regard, only one noise is considered  
 463 at a time (the other noises are deactivated in the model)  
 464 and the parameters corresponding to that noise are set based  
 465 on the values given in Table I. For the dark current and  
 466 read noise, their corresponding parameters ( $(\sigma_{\text{dark}}^i)^2$  and  
 467  $\sigma_{\text{read}}$  respectively) are incremented by 10% and the change  
 468 in the SNR value is averaged over 200 samples drawn in  
 469 each trial. In Table I, the dark current parameter is given for  
 470 the temperature of 20°C. Based on the dark current versus  
 471 temperature curve given in Ref. 25 for a CCD image sensor,  
 472 to increase the dark current by 10% at 20°C, the temperature  
 473 should go up by approximately 1°C–2°C. The read noise  
 474 parameter depends on the type of image sensor (CCD or  
 475 CMOS) and the ISO setting of the camera. In Fig. 2 of Ref. 30,  
 476 the read noise values of three image sensors are compared,  
 477 and it is indicated that changing the ISO setting of a CCD  
 478 chip between the consecutive steps may change the read noise  
 479 standard deviation by around 10–20%.

480 The SNR change can be obtained by the following  
 481 formula:

$$482 \quad \Delta\text{SNR} (\%) = 100 \times \frac{\text{SNR}_1 - \text{SNR}_2}{\text{SNR}_1}, \quad (17)$$

483 where  $\text{SNR}_1$  and  $\text{SNR}_2$  represent the SNR values before  
 484 and after incrementing the parameters respectively. Since  
 485 the noise parameters used for  $\text{SNR}_2$  are greater than  
 486 those of  $\text{SNR}_1$ , it is expected to have  $\text{SNR}_1 > \text{SNR}_2$ ,  
 487 and hence  $\Delta\text{SNR} > 0$ . A similar procedure is used to  
 488 evaluate the quantization noise by comparing the SNR of  
 489 the measurements with and without quantization noise. To  
 490 avoid saturation effects on the results, the intensity factor  
 491 is set to  $F \in \{0.1, 0.05, 0.01, 0.005, 0.001\}$ . This analysis  
 492 is performed on the boundary color patches (indices 1–13),  
 493 and the results for the data points {1, 3, 6, 8, 10, 12} together  
 494 with the average result for all 13 points are reported for the  
 495 R, G, and B sensor types in Figs 7–9. The maximum SNR  
 496 change occurs in the smallest intensity factor for the dark  
 497 current and read noise SNR sensitivity curves. However, this  
 498 pattern is not seen in the quantization noise SNR sensitivity  
 499 curves, as the R and G sensors have their maxima at different  
 500 intermediate intensities. Figure 8 shows that the SNR change  
 501 associated with read noise monotonically increases as the  
 502 light level falls off. This statement is roughly true for the  
 503 dark noise curves but it does not hold for the quantization  
 504 noise sensitivity curves. In general, no consistent pattern can  
 505 be found among the SNR sensitivity results for quantization  
 506 noise, implying that this noise does not greatly depend on  
 507 the intensity value. An interesting point that can be noted  
 508 from Figs. 7 and 8 is that for each sensor type, the data  
 509 points to which the sensor is more sensitive have lower  
 510 SNR sensitivities in comparison with other data points. For  
 511 example, in Figs. 7(a) and 8(a), the reddish color patch (index  
 512 = 6) has the least SNR sensitivity for almost all intensity

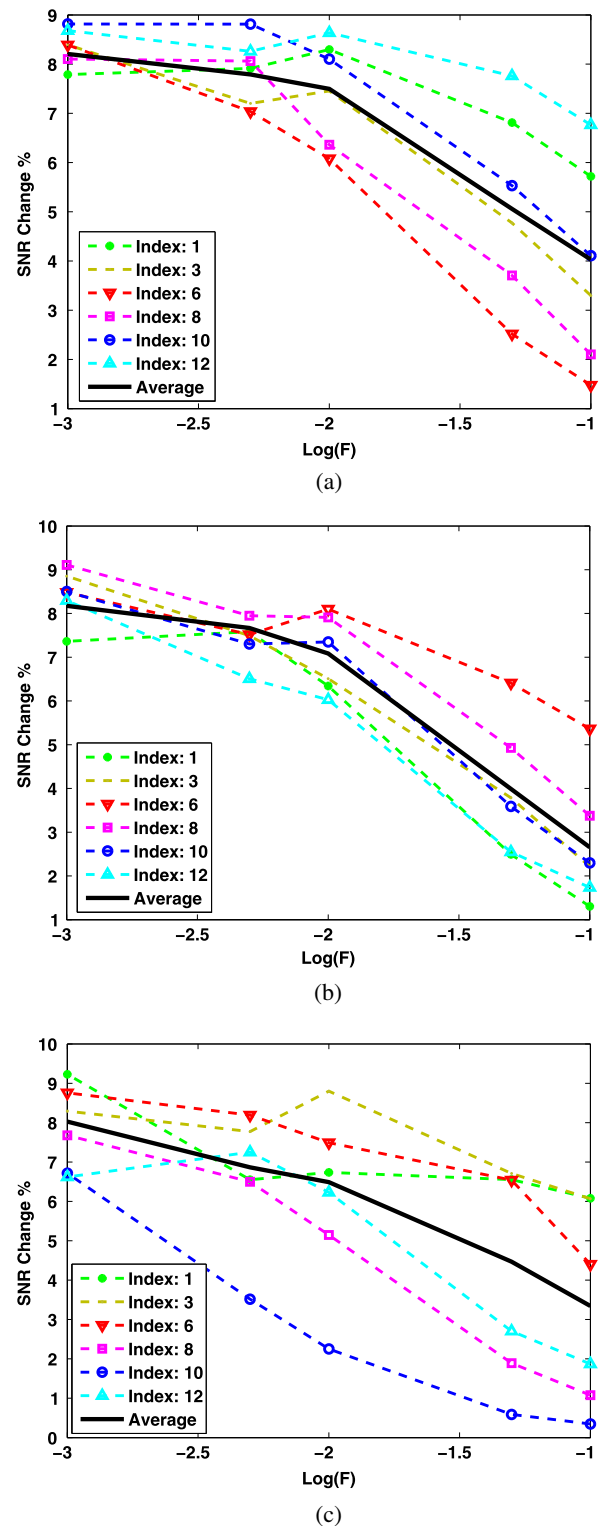
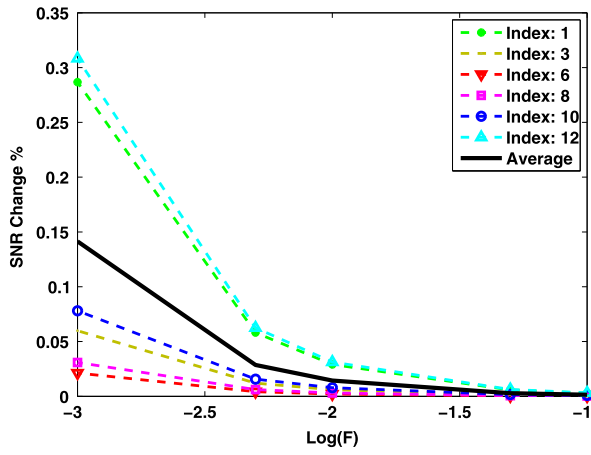
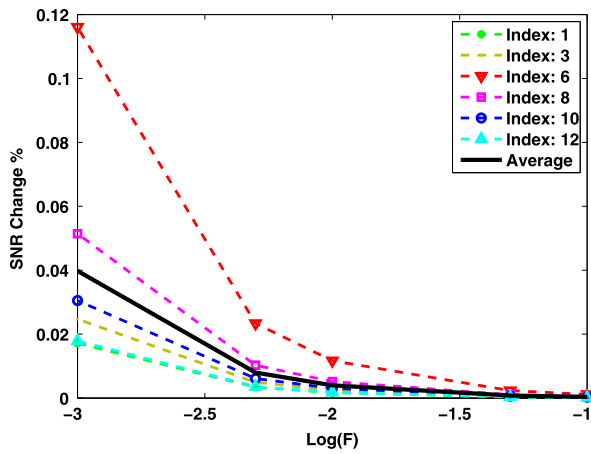


Figure 7. SNR sensitivity curves for the R, G, and B sensor types with respect to the dark current noise parameters for different color patches are plotted in (a), (b), and (c) respectively.

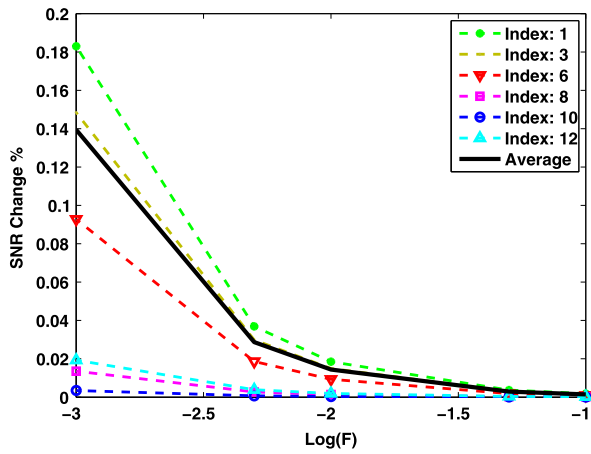
factors of the red channel; or in Figs. 7(b) and 8(b), for the green sensor, the greenish color patches (index = 1, 12) have lower SNR sensitivities compared with the other color samples. This conclusion is only true for dark current and read noise curves. Comparing the average SNR sensitivities



(a)

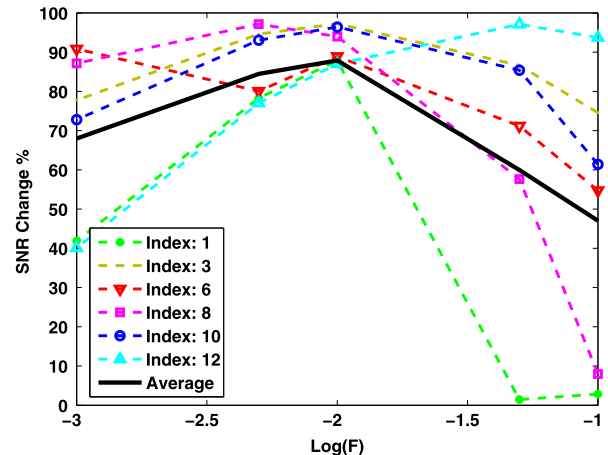


(b)

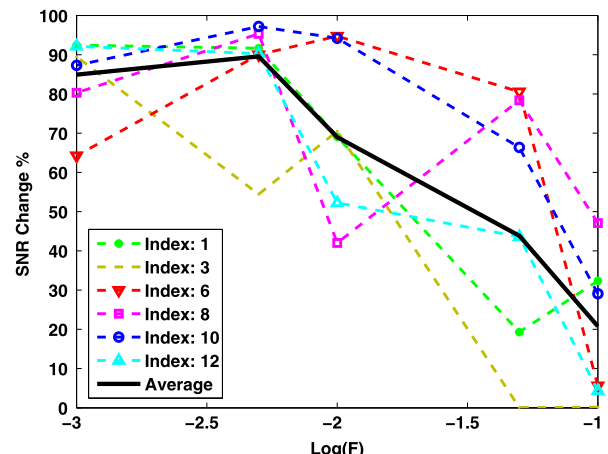


(c)

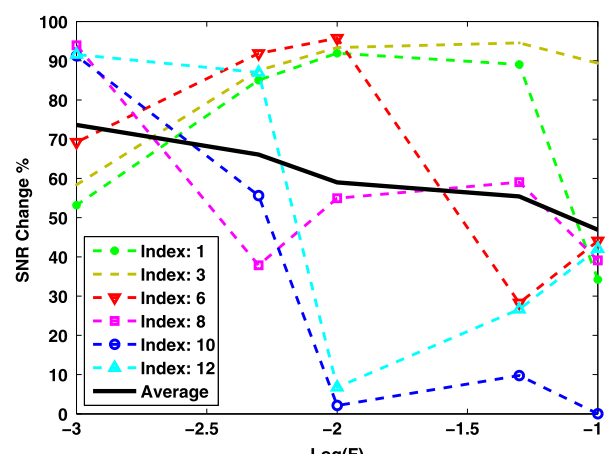
Figure 8. SNR sensitivity curves for the R, G, and B sensor types with respect to the read noise parameters for different color patches are plotted in (a), (b), and (c) respectively.



(a)



(b)



(c)

Figure 9. SNR sensitivity curves for the R, G, and B sensor types with respect to the quantization noise for different color patches are plotted in (a), (b), and (c) respectively.

518 of the three noise types reveals that read noise variations have  
 519 the least impact on the SNR (less than 1%); then dark noise  
 520 affects the SNR by between 1 and 9%; and the quantization  
 521 noise has the most significant influence on the SNR.

## CONCLUSION

522 In this work, we examined the effects of noise on color  
 523 measurements of image sensors at low light levels. In this  
 524 regard, a typical image sensor with a detailed noise model  
 525

was implemented. The image sensor model was employed in several experiments to investigate the quality of color measurements at low light intensities in the presence of noise.

The results can be summarized as follows: first, even an ideal image sensor without any device noise cannot measure consistent colors at very low light levels due to the photon noise; second, in contrast to the photon noise and read noise, which cause the generated samples at low light levels to be distributed around the high intensity samples, the dark current noise pushes the measurements towards the center of the chromaticity diagram (lower saturation values); third, for a single color patch, the distribution of noisy measured samples of an image sensor in low light situations forms an elliptical shape, and the size and inclination angle of the ellipse can give us some information about the type of noise affecting the measurements, and the color and luminance of the color patch; fourth, dark current induces a much more severe impact on color measurements in comparison to photon noise, read noise and quantization error; last but not least, the SNR sensitivity analysis showed that the presence of quantization noise does not cause a consistent change over the SNR value for different intensity factors, implying that this noise is little influenced by the light level change.

We believe that this work is relevant for many applications such as developing denoising algorithms, improving low light imaging, addressing low light image quality assessment techniques, and characterizing the noise of image sensors. Finally, this study could be further extended by incorporating the exposure time and ISO setting parameters into the model and then a set of optimal adjustments for the camera could be derived for different lighting conditions to obtain output images with the highest SNR values.

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